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## System Dynamics and Mechanical Vibration

If you have a smart project, you can say "I'm an engineer"

# Lecture 1 

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## System Dynamics and Mechanical Vibration

- Lecture aims:
- Formulate the equations of motion of two-degree-of-freedom systems
- Identify the mass, damping, and stiffness matrices from the equations of motion


## Two Degree-of-Freedom Systems

1. 

Model problem

- Matrix form of governing equation
- Special case: Undamped free vibrations
- Examples

2. Transformation of coordinates

* Inertially \& elastically coupled/uncoupled
- General approach: Modal equations
- Example

3. Response to harmonic forces

- Model equation
- Special case: Undamped system


## Two-Degree-of Freedom Systems

[ $k$ ], and [ $c$ ], [ $m$ ] are called the stiffness, damping, and mass matrices respectively,

$$
[m] \ddot{\vec{x}}(t)+[c] \dot{\vec{x}}(t)+[k] \vec{x}(t)=\vec{f}(t)
$$

$\vec{x}(t)$ and $\vec{f}(t)$ are called the displacement and force vectors, respectively.

## Two-Degree-of Freedom Systems

Thus the system has one point mass $m$ and two degrees of freedom, because the mass has two possible types of motion (translations along the $y$ and $x$ directions). The general rule for the computation of the number of degrees of freedom can be stated as follows
$\times$ number of possible types of motion of each mass


## Modeling of Mechanical System

- Spring - Mass - Damper


Force
(a)

$M \cdot \frac{d^{2}}{d t^{2}} y(t)+b \cdot \frac{d}{d t} y(t)+k \cdot y(t)=r(t)$
(b)
(a) Spring-mass-damper system.
(b) Free-body diagram.

## Two-Degree-of Freedom Systems

Equations of motion:

$$
\begin{aligned}
& m_{1} \ddot{x}_{1}(t)+\left(k_{1}+k_{2}\right) x_{1}(t)-k_{2} x_{2}(t)=0 \\
& m_{2} \ddot{x}_{2}(t)-k_{2} x_{1}(t)+\left(k_{2}+k_{3}\right) x_{2}(t)=0
\end{aligned}
$$

We are interested in knowing whether $m_{1}$ and $m_{2}$ can oscillate harmonically with the same frequency and phase angle but with different amplitudes. Assuming that it is possible to have harmonic motion of $m_{1}$ and $m_{2}$ at the same frequency $\omega$ and the same phase angle $\phi$,


## Modeling of Mechanical system

Mathematical Models for the Schematic


- Free Body Diagram FBD


## Modeling of Mechanical system

Write equation of motion: Two degree of freedom
Assume $\mathrm{X} 1>\mathrm{X} 2$ positive direction of motion $\xrightarrow{\rightarrow}$

- For mass(1)


$$
\begin{aligned}
& \sum F=-K_{d 1} x_{1}^{\prime}-K_{s 1} x_{1}-K_{d 2}\left(x_{1}{ }^{\prime}-x_{2}{ }^{\prime}\right)-K_{s 2}\left(x_{1}-x_{2}\right)=M_{1} x_{1}{ }^{\prime \prime} \\
& \quad x_{1}{ }^{\prime}\left(M_{1}\right)+x_{1}^{\prime}\left(K_{d 1}+K_{d 2}\right)+x_{1}\left(K_{s 1}+K_{s 2}\right)+x_{2}^{\prime}\left(-K_{d 2}\right)+x_{2}\left(-K_{s 2}\right)=0
\end{aligned}
$$

## Modeling of Mechanical system

Write equation of motion: Two degree of freedom
Assume X1 $>$ X2 positive direction of motion $\xrightarrow{\rightarrow}$

- For mass(2)


$$
\begin{aligned}
& \sum F=K_{d 2}\left(x_{1}{ }^{\prime}-x_{2}{ }^{\prime}\right)+K_{s 2}\left(x_{1}-x_{2}\right)+F-K_{d 3} x_{2}{ }^{\prime}-K_{s 3} x_{2}=M_{2} x_{2}{ }^{\prime \prime} \\
& \quad x_{2}{ }^{\prime \prime}\left(M_{2}\right)+x_{2}{ }^{\prime}\left(K_{d 2}+K_{d 3}\right)+x_{2}\left(K_{s 2}+K_{s 3}\right)+x_{1}{ }^{\prime}\left(-K_{d 2}\right)+x_{1}\left(-K_{s 2}\right)=F
\end{aligned}
$$

## Two-Degree-of Freedom Systems

Equations of motion:
$[m] \ddot{\vec{x}}(t)+[c] \dot{\vec{x}}(t)+[k] \vec{x}(t)=\vec{f}(t)$

(a)

$$
\begin{aligned}
& m_{1} \ddot{x}_{1}+\left(c_{1}+c_{2}\right) \dot{x}_{1}-c_{2} \dot{x}_{2}+\left(k_{1}+k_{2}\right) x_{1}-k_{2} x_{2}=f_{1} \\
& m_{2} \ddot{x}_{2}-c_{2} \dot{x}_{1}+\left(c_{2}+c_{3}\right) \dot{x}_{2}-k_{2} x_{1}+\left(k_{2}+k_{3}\right) x_{2}=f_{2}
\end{aligned}
$$



## Modeling of Mechanical System


(a) Torsional spring-mass system.
(b) Spring element.

Example

Moment of Inertia
$J^{\boldsymbol{H}}{ }^{t} \quad T=\boldsymbol{J} \ddot{\theta}$


## Two-Degree-of Freedom Systems

Equations of motion:

$$
\begin{gathered}
m_{1} \ddot{x}_{1}+\left(c_{1}+c_{2}\right) \dot{x}_{1}-c_{2} \dot{x}_{2}+\left(k_{1}+k_{2}\right) x_{1}-k_{2} x_{2}=f_{1} \\
m_{2} \ddot{x}_{2}-c_{2} \dot{x}_{1}+\left(c_{2}+c_{3}\right) \dot{x}_{2}-k_{2} x_{1}+\left(k_{2}+k_{3}\right) x_{2}=f_{2} \\
{[m] \ddot{\vec{x}}(t)+[c] \dot{\vec{x}}(t)+[k] \vec{x}(t)=\vec{f}(t)}
\end{gathered}
$$

$$
[m]=\left\lfloor\begin{array}{cc}
m_{1} & 0 \\
0 & m_{2}
\end{array}\right\rfloor
$$

$$
[c]=\left[\begin{array}{cc}
c_{1}+c_{2} & -c_{2} \\
-c_{2} & c_{2}+c_{3}
\end{array}\right]
$$


(a)

$$
[k]=\left[\begin{array}{cc}
k_{1}+k_{2} & -k_{2} \\
-k_{2} & k_{2}+k_{3}
\end{array}\right]
$$

## Two-DOF model problem

Matrix form of governing equation:


Impose the tho
dioplarements, $D^{\prime}$ Alembats $\sum F_{x}=0$


## Two-DOF model problem

Matrix form of governing equation:


Note: Matrices have positive diagonals and are symmetric.

## Undamped free vibrations

## Zero damping matrix $[\mathrm{C}]$ and force vector $\{\mathrm{P}\}$

- Equation of motion (free-undamped): $M \mathbf{x}^{\prime \prime}+K \mathbf{x}=\mathbf{0}$
- Assumed general solutions:

$$
\left\{\begin{array}{l}
x_{1} \\
x_{2}
\end{array}\right\}=\left\{\begin{array}{l}
A_{1} \\
A_{2}
\end{array}\right\} \cos (\omega t-\phi)
$$

- Differentiating twice with respect to time: (Acceleration)

$$
\{\ddot{x}\}=-\omega^{2}\left\{\begin{array}{l}
A_{1} \\
A_{2}
\end{array}\right\} \cos (\omega t-\phi)
$$

Substitute in equation of motion: (Characteristic equation)

$$
\left(K-\omega^{2} M\right)\left\{\begin{array}{l}
A_{1} \\
A_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}
$$

## Undamped free vibrations

## Zero damping matrix $[\mathrm{C}]$ and force vector $\{\mathrm{P}\}$

Characteristic equation:

$$
\left[\begin{array}{cc}
\left(k_{1}+k_{2}-m_{1} \omega^{2}\right) & -k_{2} \\
-k_{2} & \left(k_{2}-m_{2} \omega^{2}\right)
\end{array}\right]\left\{\begin{array}{l}
A_{1} \\
A_{2}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}
$$

- Characteristic polynomial (for det[ ]=0):

$$
\omega^{4}-\left(\frac{k_{1}+k_{2}}{m_{1}}+\frac{k_{2}}{m_{2}}\right) \omega^{2}+\frac{k_{1} k_{2}}{m_{1} m_{2}}=0
$$

Eigenvalues (characteristic values):

$$
\lambda_{1}=\omega_{2}^{2}=\frac{1}{2}\left\{\left[\frac{k_{1}+k_{2}}{m_{1}}+\frac{k_{2}}{m_{2}}\right] \pm\left[\left(\frac{k_{1}+k_{2}}{m_{1}}+\frac{k_{2}}{m_{2}}\right)^{2}-\frac{4 k_{1} k_{2}}{m_{1} m_{2}}\right]^{1 / 2}\right\}
$$

## Undamped free vibrations

Special case when $\mathrm{k}_{1}=\mathrm{k}_{2}=\mathrm{k}$ and $\mathrm{m}_{1}=\mathrm{m}_{2}=\mathrm{m}$
Eigenvalues and frequencies: ${\lambda_{1}}_{2}=\left\{\begin{array}{c}\omega_{1}^{2} \\ \omega_{1}^{2}\end{array}\right\}=\left\{\begin{array}{c}0.3819 \\ 2.618\end{array}\right\} \frac{\mathrm{k}}{\mathrm{m}}$

$$
\begin{aligned}
& \omega_{1}=0.618 \sqrt{\frac{k}{m}}=\text { fundamental frequency } \\
& T=\frac{2 \pi}{\omega}=\text { fundamental period }
\end{aligned}
$$

Two mode shapes (relative participation of each mass in the motion):

$1^{\text {st }}$ mode shape $\frac{A_{2}}{A_{1}}=\frac{2 k-m \omega^{2}}{k}=\frac{1.618}{1}$

## Undamped free vibrations

Two mode shapes (relative participation of each mass in the motion):
(a) First mode

(b) Second mode

## Undamped free vibrations (UFV)

Single-DOF: $\quad x(t)=C \cos \left(\omega_{n} t+\phi\right)$
For two-DOF: $\{x\}=\left\{\begin{array}{l}x_{1}(t) \\ x_{2}(t)\end{array}\right\}=C_{1}\left\{\begin{array}{l}A_{1}^{(1)} \\ A_{2}^{(1)}\end{array}\right\} \cos \left(\omega_{1} t+\phi_{1}\right)+C_{2}\left\{\begin{array}{l}A_{1}^{(2)} \\ A_{2}^{(2)}\end{array}\right\} \cos \left(\omega_{2} t+\phi_{2}\right)$
For any set of initial conditions:
$\geqslant$ We know $\{A\}^{(1)}$ and $\{A\}^{(2)}, \omega_{1}$ and $\omega_{2}$
$\diamond \quad$ Must find $C_{1}, C_{2}, \phi_{1}$, and $\phi_{2}-$ Need 4 I.C.'s

## Multi-DOF model equation

Multi-DOF systems are so similar to two-DOF.

$$
\text { Model equation: } \quad[M]\{\ddot{x}\}+[C]\{\dot{x}\}+[K]\{x\}=\{\mathrm{Q}\}
$$

We derive using:

1) Vector mechanics (Newton or D' Alembert)
2) Hamilton's principles
3) Lagrange's equations

Notes on matrices:
metric.
Kinetic energy: $\quad T=\frac{1}{2}\{\dot{x}\}^{T}[M]\{\dot{x}\}$
Strain energy in spring: $U=\frac{1}{2}\{x\}^{T}[K]\{x\}$

- $[\mathrm{M}]$ is positive definite (since $T$ is always positive)
- $[\mathrm{K}]$ is positive semi-definite:
- all positive eigenvalues, except for some potentially 0 -eigenvalues which occur during a rigid-body motion.
- If restrained/tied down $\Rightarrow$ positive-definite. All positive.


## Inverted Pendulum



## Projects



## Projects



## Projects



## Program



Express Vibration Lab.vi

